

## ON THE GROWTH OF THE SPECTRAL MEASURE

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**ABSTRACT.** We are concerned with the asymptotics of the spectral measure associated with a self-adjoint operator. By using comparison techniques we shall show that the eigenfunctionals of  $L_2$  are close to the eigenfunctionals  $L_1$  if and only if  $d\Gamma_1 \asymp d\Gamma_2$  as  $\lambda \rightarrow \infty$ .

**KEY WORDS AND PHRASES:** Spectral asymptotics, spectral function, Sturm-Liouville operators.

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### 1 INTRODUCTION

We would like to obtain a relation between the growth of the spectral measure of a self-adjoint operator and the behaviour of its eigenfunctionals. In this study we shall assume that we have two "close" self-adjoint operators acting in the same separable Hilbert space,  $H$  say. Without loss of generality we can assume that both operators have simple spectra. To this end, let us denote by  $\varphi(\lambda)$  and  $y(\lambda)$  the eigenfunctionals of  $L_1$  and  $L_2$  respectively. Recall that the spectrum of a self-adjoint operator is defined by

$$\forall \lambda \in \sigma_i \exists \varphi_{i,n} \in D_{\{L_i\}} / \|\varphi_{i,n}\| = 1 \text{ and } \|L_i \varphi_{i,n} - \lambda \varphi_{i,n}\| \xrightarrow{n \rightarrow \infty} 0$$

where  $i = 1, 2$ . In case  $\lambda$  is in the continuous spectrum the sequence is not compact in the Hilbert space  $H$ . For this we can assume the existence of a countably normed perfect space  $\Phi$ , such that

$$\Phi \hookrightarrow H \hookrightarrow \Phi'$$

where the embeddings are compact, for further details see [1] and [2]. For the sake of simplicity we shall assume that the embeddings are given by the identities and so

$$f \in \Phi \quad \psi \in H \quad (f, \psi) \equiv \langle f, \psi \rangle_{\Phi \times \Phi'}$$

Since the sequence  $\varphi_n$  is bounded in  $H$ , it is then compact in  $\Phi'$ , which implies

$$\varphi_n \xrightarrow{\Phi'} \varphi(\lambda) \in \Phi'$$

and similarly for the operator  $L_2$ ; Since both operators are acting in the same Hilbert space  $H$ , we shall assume that the space  $\Phi'$  contains both systems of eigenfunctionals; i.e.,

$$\{y(\lambda)\} \subset \Phi' \quad \text{and} \quad \{\varphi(\lambda)\} \subset \Phi'$$

Recall that the system  $\{y(\lambda)\}$  helps define an isometry for  $L_2$

$$\forall f \in \Phi \quad f \longrightarrow \hat{f}^2(\lambda) \equiv \langle f, y(\lambda) \rangle_{\Phi \times \Phi'}$$

$$f = \int \overline{\hat{f}^2(\lambda)} y(\lambda) d\Gamma_2(\lambda) \quad \text{where} \quad \hat{f}^2(\lambda) \in L^2_{d\Gamma_2(\lambda)}$$

Similarly for  $\varphi(\lambda)$ ;

$$\forall f \in \Phi \quad f \longrightarrow \hat{f}^1(\lambda) \equiv \langle f, \varphi(\lambda) \rangle_{\Phi \times \Phi'},$$

$$f = \int \overline{\hat{f}^1(\lambda)} \varphi(\lambda) d\Gamma_1(\lambda) \quad \text{where } \hat{f}^1(\lambda) \in L^2_{d\Gamma_1(\lambda)}$$

These transforms define isometries, and Parseval equality yields

$$\int_{\sigma_1} \hat{f}^1(\lambda) \overline{\hat{\psi}^1(\lambda)} d\Gamma_1(\lambda) = (f, \psi)_H = \int_{\sigma_2} \hat{f}^2(\lambda) \overline{\hat{\psi}^2(\lambda)} d\Gamma_2(\lambda).$$

where the nondecreasing functions  $\Gamma_1(\lambda)$  and  $\Gamma_2(\lambda)$  are called the spectral measures associated with  $L_1$  and  $L_2$ , respectively. It is these functions that we would like to estimate as  $\lambda \rightarrow \infty$ .

In all that follows  $y(\lambda) \sim \varphi(\lambda)$  as  $\lambda \rightarrow \infty$  means  $\forall f \in \Phi$ ,

$$\hat{f}^1(\lambda) \asymp \hat{f}^2(\lambda) \quad \text{as } \lambda \rightarrow \infty$$

and  $d\Gamma_1(\lambda) \sim d\Gamma_2(\lambda)$  as  $\lambda \rightarrow \infty$  means that  $\forall F \in L^1_{d\Gamma_1(\lambda)} \cap L^1_{d\Gamma_2(\lambda)}$

$$\int_{\lambda}^{\infty} F(\eta) d\Gamma_1(\eta) \asymp \int_{\lambda}^{\infty} F(\eta) d\Gamma_2(\eta) \quad \text{as } \lambda \rightarrow \infty.$$

In this work, we shall try to answer the following problem:

**Statement of the Problem:** under what conditions

$$y(\lambda) \sim \varphi(\lambda) \text{ as } \lambda \rightarrow \infty \Leftrightarrow d\Gamma_1(\lambda) \sim d\Gamma_2(\lambda) \text{ as } \lambda \rightarrow \infty.$$

In order to answer the above question, we shall compare the self-adjoint operators  $L_1$  and  $L_2$ , see [3]. Recall that a shift operator or transmutation is defined by

$$y(\lambda) = V\varphi(\lambda) \quad \lambda \in \sigma_1;$$

Clearly the definition of  $V$  depends on  $\sigma_2$  and  $\sigma_1$  and we shall agree to set

$$y(\lambda) = 0 \text{ if } \lambda \notin \sigma_2, \text{ and } \varphi(\lambda) = 0 \text{ if } \lambda \notin \sigma_1$$

$$y(\lambda) = V\varphi(\lambda) \quad \lambda \in \sigma_2 \subset \sigma_1 \subset R.$$

Condition  $\sigma_2 \subset \sigma_1$  insures that  $V0 = 0$  and so defines an operator on the algebraic span of  $\{\varphi(\lambda)\}$ . Thus it is clear that in order for  $V$  and  $V^{-1}$  to exist as linear operator it is necessary that  $\sigma_2 \subset \sigma_1$  and  $\sigma_1 \subset \sigma_2$

$$\sigma_2 \equiv \sigma_1.$$

It is readily seen that  $\{\varphi(\lambda)\}$  form a complete set in the reflexive space (perfect)  $\Phi'$ , and so the space generated by  $\{\varphi(\lambda)\}$  is dense in  $\Phi'$ . Consequently  $V$  is densely defined. This in turns allows us to define the adjoint operator  $V' : \Phi \rightarrow \Phi$ .

## 2 MAIN RESULTS

We shall agree to say  $\Gamma_1(\lambda)$  is Abs- $d\Gamma_2$  if there exists  $g(\eta) \in L^{1,loc}_{d\Gamma_2}$  such that

$$\Gamma_1(\lambda) = \int_0^{\lambda} g(\eta) d\Gamma_2(\eta) + \Gamma_1(0)$$

This fact shall be denoted by

$$g(\lambda) \equiv \frac{d\Gamma_1}{d\Gamma_2}(\lambda) \in L^{1,loc}_{d\Gamma_2}$$

In this case the condition  $d\Gamma_1(\lambda) \sim d\Gamma_2(\lambda)$  in the statement of the problem can be restated as  $g(\lambda) \asymp 1$  as  $\lambda \rightarrow \infty$ . Recall that due to reflexivity of the space  $\Phi$ , the operator  $V'$  is defined in  $\Phi$  and since  $\Phi \hookrightarrow H$ ,  $V'$  is actually defined in  $H$ . Let us denote this extension to the space  $H$  by  $\tilde{V}$ . Since we are interested in the case where  $y(\lambda) \sim \varphi(\lambda)$  we can expect  $V$  to be bounded. In this regard we have the following result:

**Theorem 1:** If the extension  $\tilde{V} : H \rightarrow H$ , is a bounded operator then  $\Gamma_1(\lambda)$  is  $d\Gamma_2$ -ABS continuous.

**Proof:** It is clear that for  $f \in D_{V'}$

$$\begin{aligned} \langle f, y(\lambda) \rangle_{\Phi \times \Phi'} &= \langle f, V\varphi(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \langle V'f, \varphi(\lambda) \rangle_{\Phi \times \Phi'} \end{aligned}$$

In other words

$$\hat{f}^2(\lambda) = \widehat{V'f}^1(\lambda). \tag{2.1}$$

Equation 2.1 obviously holds for  $f \in H$ . Indeed let  $f_n \in D_{V'} \subset H$  such that  $f_n \xrightarrow{H} f \in H$ . Given that  $\tilde{V}$  is a bounded operator in  $H$ , we obviously have  $\tilde{V}f_n \rightarrow \tilde{V}f$ . Using the fact that  $\forall n, \hat{f}_n^2(\lambda) = \widehat{V'f_n}^1(\lambda)$  and the isometries are bounded operators we have  $\hat{f}_n^2 \rightarrow \hat{f}^2$  and  $\widehat{V'f_n}^1 \rightarrow \widehat{V'f}^1$ . Therefore

$$\hat{f}^2(\lambda) = \widehat{V'f}^1(\lambda) \quad f \in H. \tag{2.2}$$

From which we deduce that  $\forall f \in H$

$$\begin{aligned} \int \hat{f}^2(\lambda) \overline{\hat{f}^2(\lambda)} d\Gamma_1(\lambda) &= \int \widehat{V'f}^1 \overline{\widehat{V'f}^1} d\Gamma_1(\lambda) \\ &= (\tilde{V}'f, \tilde{V}'f) \\ &= \|\tilde{V}'f\|^2 \\ &\leq c\|f\|^2 \\ &\leq c \int |\hat{f}^2(\lambda)|^2 d\Gamma_2(\lambda) \quad \forall f \in H. \end{aligned}$$

Thus each  $d\Gamma_2$  negligible set is a  $d\Gamma_1$  negligible set. Henceforth  $\Gamma_1(\lambda)$  to be  $d\Gamma_2(\lambda)$ -Abs continuous. The above inequality is exactly a sufficient condition for the Radon-Nikodym theorem to hold, see [4].

In all that follows we shall assume that  $d\Gamma_1(\lambda)$  is  $d\Gamma_2 - Abs$  continuous which is denoted by

$$g(\lambda) \equiv \frac{d\Gamma_1}{d\Gamma_2}(\lambda).$$

We now need to define a function of an operator, namely  $g(L_2)$  for the next result:

$$\begin{aligned} \Phi &\xrightarrow{g(L_2)} H \\ f &\longrightarrow g(L_2)f \equiv \int g(\lambda) \overline{\hat{f}^2(\lambda)} y(\lambda) d\Gamma_2(\lambda). \end{aligned}$$

**Theorem 2:** Assume that  $V$  admits closure in  $\Phi'$  and  $\Gamma_1$  is Abs- $d\Gamma_2(\lambda)$  then

$$\forall \psi \in D_{V'} \subset \Phi \quad \left( \sqrt{\frac{d\Gamma_2}{d\Gamma_1}(L_2)} \right)' \left( \sqrt{\frac{d\Gamma_2}{d\Gamma_1}(L_2)} \right) \psi = \overline{V}V'\psi \quad \text{in } \Phi'.$$

**Proof:** From equation 2.1 and the fact that the embeddings are defined by identities, we deduce that  $\forall f, \psi \in D_{V'} \subset \Phi$

$$\int \hat{f}^2(\lambda) \overline{\hat{\psi}^2(\lambda)} d\Gamma_1(\lambda) = \int \widehat{V'f}^1 \overline{\widehat{V'\psi}^1} d\Gamma_1(\lambda) \tag{2.3}$$

$$\begin{aligned} &= (V'f, V'\psi) \tag{2.4} \\ &= \langle V'f, V'\psi \rangle_{\Phi \times \Phi'} \\ &= \langle f, \overline{V}V'\psi \rangle_{\Phi \times \Phi'}. \end{aligned}$$

However the left handside of equation 2.3 can rewritten as

$$\begin{aligned}
 \int \hat{f}^2(\lambda) \overline{\hat{\psi}^2(\lambda)} d\Gamma_1(\lambda) &= \int \hat{f}^2(\lambda) \overline{\hat{\psi}^2(\lambda)} g(\lambda) d\Gamma_2(\lambda) \\
 &= \int \sqrt{g(\lambda)} \hat{f}^2(\lambda) \sqrt{g(\lambda)} \overline{\hat{\psi}^2(\lambda)} d\Gamma_2(\lambda) \\
 &= \int \widehat{\sqrt{g(L_2)} f} \widehat{\sqrt{g(L_2)} \psi}^2 d\Gamma_2(\lambda) \\
 &= (\sqrt{g(L_2)} f, \sqrt{g(L_2)} \psi) \\
 &= \langle f, \sqrt{g(L_2)}' \sqrt{g(L_2)} \psi \rangle_{\Phi \times \Phi'}.
 \end{aligned}
 \tag{2.5}$$

Observe that if we set  $f = \psi$  in equations 2.4 and 2.5 then we would obtain

$$\|\sqrt{g(L_2)} f\| = \|V' f\|
 \tag{2.6}$$

from which we deduce that  $D_{V'} \subset D_{\sqrt{g(L_2)}} \subset \Phi$ , from we obtain

$$\forall \psi \in D_{V'}, \quad \sqrt{g(L_2)}' \sqrt{g(L_2)} \psi = \overline{V} V' \psi.
 \tag{2.7}$$

**Remark:** Observe that both operators  $\sqrt{g(L_2)}' \sqrt{g(L_2)}$  and  $\overline{V} V'$  are mappings from  $\Phi \longrightarrow \Phi'$ .

It is easy to see that if we restrict equation 2.7 to

$$f \in D_{g(L_2)} \equiv \{f \in \Phi / g(\lambda) \hat{f}^2(\lambda) \in L^2_{d\Gamma_2}\}$$

then it reduces to

$$\forall f \in D_{V'} \cap D_{g(L_2)} \quad \frac{d\Gamma_1}{d\Gamma_2}(L_2) = g(L_2) = \overline{V} V' \quad \text{in } \Phi'
 \tag{2.8}$$

The next result describes the domain of  $\tilde{V}'$ .

**Theorem 3:**  $\tilde{V}$  is densely defined if and only if  $L^2_{d\Gamma_1(\lambda)} \cap L^2_{d\Gamma_2(\lambda)}$  is dense in  $L^2_{d\Gamma_2(\lambda)}$ .

**Proof:** From equation 2.2 it is readily seen that

$$f \in D_{\tilde{V}} \Leftrightarrow \hat{f}^2(\lambda) \in L^2_{d\Gamma_1(\lambda)} \cap L^2_{d\Gamma_2(\lambda)}$$

Then use the fact that  $f \longrightarrow \hat{f}^2$  is an isometry between  $H$  and  $L^2_{d\Gamma_2(\lambda)}$ .

This work is based on the following result.

**Theorem 4:** Assume that

- $V$  admits closure in  $\Phi'$
- $\Gamma_1$  is Abs-  $d\Gamma_2(\lambda)$
- $\overline{V}^{-1}$  exists
- $\overline{V} : \Phi \longrightarrow \Phi$  is a bounded operator

then

$$g(\lambda) \varphi(\lambda) - y(\lambda) = (V' - 1)y(\lambda) \quad \text{in } \Phi'.$$

**Proof:** Notice that conditions of Theorem 2 hold and so it follows that

$$\sqrt{g(L_2)}' \sqrt{g(L_2)} = \overline{V} V' \quad \text{in } \Phi'.
 \tag{2.9}$$

By the above condition we have that  $\sqrt{g(L_2)}' \sqrt{g(L_2)} f \in \Phi$  if  $f \in D_{V'} \subset \Phi$ . However since it is assumed that  $\overline{V}^{-1}$  exists, then equation 2.8 yields

$$\overline{V}^{-1} \left( \sqrt{g(L_2)}' \right)' \sqrt{g(L_2)} = V' \quad \text{in } \Phi'
 \tag{2.10}$$

In order to proceed further we need to extend the operator  $V'$  to  $\Phi'$ . For this observe that since  $\overline{V} : \Phi \longrightarrow \Phi$  is a bounded operator,  $\overline{V}' = V'$  is a bounded operator in  $\Phi'$ . Hence  $V'$  is defined for all elements in  $\Phi'$ , and in particular for  $y(\lambda)$ , thus

$$V^{-1}\sqrt{g(L_2)'}\sqrt{g(L_2)}y(\lambda) = V'y(\lambda).$$

We now need to compute  $\sqrt{g(L_2)'}\sqrt{g(L_2)}y(\lambda)$ . Let  $f \in D_{V'} \subset \Phi$  then

$$\begin{aligned} \langle f, \sqrt{g(L_2)'}\sqrt{g(L_2)}y(\lambda) \rangle_{\Phi \times \Phi'} &= \langle \sqrt{g(L_2)'}f, \sqrt{g(L_2)}y(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \langle \sqrt{g(L_2)'}\sqrt{g(L_2)}f, y(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \frac{1}{\sqrt{g(\lambda)}\sqrt{g(\lambda)}}f^2(\lambda) \\ &= g(\lambda)f^2(\lambda) \\ &= \langle f, g(\lambda)y(\lambda) \rangle \end{aligned}$$

where we have used the fact that  $\sqrt{g(L_2)'}\sqrt{g(L_2)}f = \bar{V}V'f \in \Phi$ . Hence

$$\sqrt{g(L_2)'}\sqrt{g(L_2)}y(\lambda) = g(\lambda)y(\lambda) \text{ in } \Phi' \text{ d}\Gamma_2 \text{ a.e.}$$

where  $g(\lambda) \equiv \frac{d\Gamma_1}{d\Gamma_2}(\lambda)$  is a real function. Hence we have

$$g(\lambda)\bar{V}^{-1}y(\lambda) = V'y(\lambda).$$

Since by definition we have  $\bar{V}^{-1}y(\lambda) = \varphi(\lambda)$  we obtain

$$g(\lambda)\varphi(\lambda) - y(\lambda) = (V' - 1)y(\lambda) \text{ in } \Phi'.$$

We easily deduce the following result:

**Corollary 1:** Let conditions of Theorem 4 hold then

$$g(\lambda)\varphi(\lambda) - y(\lambda) \xrightarrow{\Phi'} 0 \Leftrightarrow (V' - 1)y(\lambda) \xrightarrow{\Phi'} 0$$

**Corollary 2:** Let conditions of Theorem 4 hold and  $(V' - 1)y(\lambda) \xrightarrow{\Phi'} 0$   $\lambda \rightarrow \infty$  then

$$g(\lambda) \sim 1 \text{ as } \lambda \rightarrow \infty \Leftrightarrow \varphi(\lambda) - y(\lambda) \xrightarrow{\Phi'} 0 \text{ as } \lambda \rightarrow \infty.$$

**Proof:** By hypothesis and Corollary 1 we have  $\forall f \in \Phi$

$$g(\lambda)\hat{f}^1(\lambda) - \hat{f}^2(\lambda) \rightarrow 0 \text{ as } \lambda \rightarrow \infty.$$

Thus if  $g(\lambda) \rightarrow 1$  then  $\hat{f}^1(\lambda) - \hat{f}^2(\lambda) \rightarrow 0$  which means that  $\varphi(\lambda) - y(\lambda) \xrightarrow{\Phi'} 0$  as  $\lambda \rightarrow \infty$ .

Conversely  $\hat{f}^1(\lambda) - \hat{f}^2(\lambda) \rightarrow 0$  together with  $y(\lambda) - g(\lambda)\varphi(\lambda) \xrightarrow{\Phi'} 0$  implies that

$$g(\lambda)\hat{f}^1(\lambda) - \hat{f}^2(\lambda) \rightarrow 0$$

i.e.  $g(\lambda) \rightarrow 1$  as  $\lambda \rightarrow \infty$ .

Corollary 2 suggests to write  $V = 1 + K$ . In this case Theorem 2 would read

$$g(\lambda)\varphi(\lambda) - y(\lambda) \xrightarrow{\Phi'} 0 \Leftrightarrow K'y(\lambda) \xrightarrow{\Phi'} 0 \text{ as } \lambda \rightarrow \infty.$$

The question we would like to answer now is under what condition would

$$K'y(\lambda) \xrightarrow{\Phi'} 0 \text{ as } \lambda \rightarrow \infty.$$

First we need to observe that the above convergence holds in  $\Phi'$ . Indeed by construction the function  $y(\lambda)$  is in  $\Phi'$  and so the operator  $K'$  originally was defined in  $\Phi$  must be extended to  $\Phi'$ .

This is easily achieved if the operator  $K$ , i.e.  $\bar{V}$ , is bounded in  $\Phi \rightarrow \Phi$ .

**Theorem 5:** Let

- $V : \Phi \rightarrow \Phi$  be a bounded operator.
- $K \equiv \bar{V} - 1$ , be such that  $\Phi \xrightarrow{L_2K} H$  is densely defined in  $\Phi$

then

$$K'y(\lambda) \xrightarrow{\Phi'} 0 \text{ as } \lambda \rightarrow \infty.$$

**Proof:** Recall that for each  $\lambda$ , there exists a bounded sequence  $\varphi_{n,\lambda} \in D_{L_2}$  such that

$$\varphi_{n,\lambda} \in D_{L_2}, \quad \|\varphi_{n,\lambda}\| = 1, \quad \text{and} \quad \|L_2\varphi_{n,\lambda} - \lambda\varphi_{n,\lambda}\| \longrightarrow 0$$

The last condition can be written as

$$\lambda\varphi_{n,\lambda} = L_2\varphi_{n,\lambda} + \epsilon(n, \lambda)$$

where  $\epsilon(n, \lambda) \rightarrow 0$  in  $H$  as  $n \rightarrow \infty$ . This allows us to obtain the following limit

$$\begin{aligned} \langle f, K'y(\lambda) \rangle_{\Phi \times \Phi'} &= \langle Kf, y(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \lim_{n \rightarrow \infty} \langle Kf, \varphi_{n,\lambda} \rangle_{\Phi \times \Phi'} \\ &= \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle \lambda\varphi_{n,\lambda}, Kf \rangle \\ &= \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle L_2\varphi_{n,\lambda} + \epsilon(n, \lambda), Kf \rangle \\ &= \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle L_2\varphi_{n,\lambda}, Kf \rangle + \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle \epsilon(n, \lambda), Kf \rangle \\ &= \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle \varphi_{n,\lambda}, L_2Kf \rangle + \frac{1}{\lambda} \lim_{n \rightarrow \infty} \langle \epsilon(n, \lambda), Kf \rangle \\ &\leq \frac{1}{\lambda} \|\varphi_{n,\lambda}\| \|L_2Kf\| + \frac{1}{\lambda} \lim_{n \rightarrow \infty} \|\epsilon(n, \lambda)\| \|Kf\| \end{aligned}$$

So as  $\lambda \rightarrow \infty$  we shall obtain  $\langle f, K'y(\lambda) \rangle_{\Phi \times \Phi'} \rightarrow 0$ . This last limit means that

$$K'y(\lambda) \xrightarrow{\Phi'} 0 \quad \text{as} \quad \lambda \rightarrow \infty.$$

Recall that in order for the conclusion to hold we need  $L_2K$  to be at least densely defined in  $\Phi$ .

**Remark:** The condition  $V : \Phi \rightarrow \Phi$  bounded can be replaced by densely defined. This forces us to use Baire's Theorem to obtain the density of  $\Phi \cap D_V \cap D_{L_2K}$  in  $\Phi$ .

**Theorem 6:** Let the conditions of Theorem 2 hold, and

- $V : \Phi \rightarrow \Phi$  be a bounded operator
- $(g(L_2) - 1)^{-1}K$  be a bounded operator in  $\Phi$

then

$$(g(\lambda) - 1)y(\lambda) \xrightarrow{\Phi'} 0 \quad \Rightarrow \quad K'y(\lambda) \rightarrow 0 \quad \text{as} \quad \lambda \rightarrow \infty.$$

**Proof:**

$$\begin{aligned} \langle f, K'y(\lambda) \rangle_{\Phi \times \Phi'} &= \langle Kf, y(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \widehat{Kf^2}(\lambda) \\ &= (g(\lambda) - 1)(g(\lambda) - 1)^{-1}\widehat{Kf^2}(\lambda) \\ &= (g(\lambda) - 1)\{(g(L_2) - 1)^{-1}Kf\}^2 \\ &= (g(\lambda) - 1) \langle (g(L_2) - 1)^{-1}Kf, y(\lambda) \rangle_{\Phi \times \Phi'} \\ &= \langle (g(L_2) - 1)^{-1}Kf, (g(\lambda) - 1)y(\lambda) \rangle_{\Phi \times \Phi'} \end{aligned}$$

Since the  $[g(\lambda) - 1]y(\lambda) \xrightarrow{\Phi'} 0$  we obtain  $\langle f, K'y(\lambda) \rangle_{\Phi \times \Phi'} \rightarrow 0 \quad \forall f \in \Phi$  and so  $K'y(\lambda) \rightarrow 0$  as  $\lambda \rightarrow \infty$ .

**Corollary 3:** Assume that conditions of Theorem 4, hold and

- $y(\lambda)$  are bounded functionals for large  $\lambda$
- $(g(L_2) - 1)^{-1}K$  be a bounded operator in  $\Phi$

then

$$g(\lambda) - 1 \xrightarrow{\lambda \rightarrow \infty} 0 \Rightarrow y(\lambda) - \varphi(\lambda) \xrightarrow{\Phi'} 0 \text{ as } \lambda \rightarrow \infty$$

**Proof:** It suffices to see that  $(g(\lambda) - 1)y(\lambda) \xrightarrow{\Phi'} 0$ , and since Theorem 6, is applicable

$$K'y(\lambda) \rightarrow 0 \text{ as } \lambda \rightarrow \infty.$$

From Theorem 4, we deduce that

$$g(\lambda)\varphi(\lambda) - y(\lambda) \xrightarrow{\Phi'} 0.$$

It remains to see that since  $g(\lambda) \asymp 1$  as  $\lambda \rightarrow \infty \Rightarrow \varphi(\lambda) \xrightarrow{\Phi'} y(\lambda)$  as  $\lambda \rightarrow \infty$ .

### 3 EXAMPLES

Below we shall consider two simple examples to illustrate the above results.

Let  $L_1$  and  $L_2$  be two self-adjoint differential operators in  $L^2[0, \infty)$  defined by

$$\begin{cases} L_1 f \equiv -f''(x) + q(x)f(x) \\ n f(0) - f'(0) = 0 \end{cases} \quad \text{and} \quad \begin{cases} L_2 f \equiv -f''(x) \\ n f(0) - f'(0) = 0. \end{cases}$$

where  $|n| < \infty$ . Let the eigenfunctionals associated with  $L_1$  and  $L_2$  be defined by

$$\begin{cases} L_1 \varphi(x, \lambda) \equiv \lambda \varphi(x, \lambda) \\ \varphi(0, \lambda) = 1, \quad \varphi'(0, \lambda) = n \end{cases} \quad \text{and} \quad \begin{cases} L_2 y(x, \lambda) = \lambda y(x, \lambda) \\ y(0, \lambda) = 1, \quad y'(0, \lambda) = n \end{cases}$$

where  $y(x, \lambda) = \cos(\sqrt{\lambda}x) + n \frac{\sin(\sqrt{\lambda}x)}{\sqrt{\lambda}}$ . It is clear that

$$\varphi(x, \lambda) = y(x, \lambda) + \int_0^x \frac{\sin(\sqrt{\lambda}(x-t))}{\sqrt{\lambda}} q(t) \varphi(t, \lambda) dt.$$

By the Riemman-Lebesgue theorem we have

$$\varphi(x, \lambda) - y(x, \lambda) \rightarrow 0 \text{ as } \lambda \rightarrow \infty.$$

It is also known that the following representation holds

$$\varphi(x, \lambda) = y(x, \lambda) + \int_0^x K(x, t)y(t, \lambda) dt.$$

Then formally

$$(V' - 1)y(x, \lambda) = \int_x^\infty K(t, x)y(t, \lambda) dt$$

Therefore if  $(V' - 1)y(x, \lambda) \xrightarrow{\Phi'} 0$  then

$$\frac{d\Gamma_1(\lambda)}{d\lambda} \asymp \frac{1}{\pi} \frac{\sqrt{\lambda}}{\lambda + n^2} \text{ as } \lambda \rightarrow \infty.$$

**Remark:** It is known that if  $q'(x) \in L^{1,loc}[0, \infty)$  then for each fixed  $x$   $K_{tt}(x, t) \in L^{1,loc}[0, \infty)$  and hence  $L_2 K$  is densely defined. Therefore Theorem 5 is applicable.

The next example deals with the generalized Sturm Liouville operator. Let

$$\begin{cases} L_1 f \equiv -\frac{1}{w(x)} f''(x) + q(x)f(x) \\ f'(0) = 0. \end{cases} \quad \text{and} \quad \begin{cases} L_2 f \equiv \frac{-1}{x^\alpha} f'(x) \\ f'(0) = 0. \end{cases}$$

where  $w(x) \asymp x^\alpha$  as  $x \rightarrow 0$  and  $\alpha > 0$ . In this case the operator  $L_2$  corresponds to a string whose length and mass are infinite, and is known to be self-adjoint in the space  $L^2_{x^\alpha dx}$ , see [5, p. 151] and [9].

We shall see that the behaviour of  $w(x) \rightarrow 0$  dictates the behaviour of the spectral function at infinity. Although this result is known, see [6], we shall provide a different treatment as it is stated in [7]. For simplicity let the eigenfunctionals associated with  $L_1$  and  $L_2$  be defined by

$$\begin{cases} L_1\varphi(x, \lambda) \equiv \lambda\varphi(x, \lambda) \\ \varphi(0, \lambda) = 1, \varphi'(0, \lambda) = 0 \end{cases} \quad \text{and} \quad \begin{cases} L_2y(x, \lambda) = \lambda y(x, \lambda) \\ y(0, \lambda) = 1, y'(0, \lambda) = 0. \end{cases}$$

It is clear that

$$\varphi(x, \lambda) = y(x, \lambda) + \int_0^x R(x, t, \lambda)q(t)\varphi(t, \lambda)dt.$$

where  $R(x, t, \lambda)$  is the Greens' function and it is shown, by the semi-classical approximation, see [8], that  $R(x, t, \lambda) \rightarrow 0$  as  $\lambda \rightarrow \infty$ . Therefore we have that  $\varphi(x, \lambda) - y(x, \lambda) \rightarrow 0$  as  $\lambda \rightarrow \infty$ . The solution  $y(x, \lambda)$  are known explicitly,

$$y(x, \lambda) = \sqrt{x}AJ_{-\nu}\left(\frac{2\sqrt{\lambda}}{\alpha+2}x^{\frac{\alpha+2}{2}}\right).$$

where  $\nu = \frac{1}{\alpha+2}$  and  $A = \left\{\frac{2\sqrt{\lambda}}{\alpha+2}\right\}^{\frac{1}{\alpha+2}} \frac{1}{\Gamma(1-\nu)}$ .

Therefore provided  $(V' - 1)y(x, \lambda) \xrightarrow{\Phi'} 0$ , we shall have

$$\Gamma_1(\lambda) \asymp \Gamma_2(\lambda) \quad \text{as } \lambda \rightarrow \infty.$$

where, see [3],  $\Gamma_2(\lambda) = c\lambda^{\frac{\alpha+1}{\alpha+2}}$  for  $\lambda > 0$ .

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