Thermodynamical Limit for Correlated Gaussian Random Energy Models

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Abstract

Let $\{E_\sigma(N)\}_{\sigma\in\Sigma_N}$ be a family of $|\Sigma_N| = 2^N$ centered unit Gaussian random variables defined by the covariance matrix C_N of elements $c_N(\sigma,\tau) := \mathrm{Av}\,(E_\sigma(N)E_\tau(N))$, and $H_N(\sigma) = -\sqrt{N}E_\sigma(N)$ the corresponding random Hamiltonian. Then the quenched thermodynamical limit exists if, for every decomposition $N = N_1 + N_2$, and all pairs $(\sigma,\tau) \in \Sigma_N \times \Sigma_N$:

$$c_N(\sigma,\tau) \le \frac{N_1}{N}\, c_{N_1}(\pi_1(\sigma),\pi_1(\tau)) + \frac{N_2}{N}\, c_{N_2}(\pi_2(\sigma),\pi_2(\tau))$$

where $\pi_k(\sigma), k = 1,2$ are the projections of $\sigma \in \Sigma_N$ into Σ_{N_k}. The condition is explicitly verified for the Sherrington-Kirckpatrick, the even p-spin, the Derrida REM and the Derrida-Gardner GREM models.

1 Introduction, Definitions and Results

It has recently been proved by Guerra and Toninelli [GuTo] that for the Sherrington-Kirckpatrick (hereafter SK) model (as well as for the even-p-spin models) the thermodynamical limit exists for the quenched free energy and almost everywhere for its random realizations. In this paper we single out general sufficient conditions that imply the existence of the quenched thermodynamical limit for any correlated Gaussian random energy model. Our analysis thus includes as special cases not only the even p spin models (in particular the SK one, $p = 2$) but also the Derrida REM model[De1],[De2] and the Derrida-Gardner GREM[DeGa].

The paper is organized as follows: in this section we introduce the definitions and state the results. In section 3, after introducing and elucidating the operation of *lifting* for a family of Gaussian random variables, we describe the proof of our theorem. In section 4 we show how our analysis can be applied to the specific examples listed above.

To define the set up we consider a disordered model having 2^N energy levels where N is the size of the system. We label the energy levels by the index $\sigma = \{\sigma_1, \sigma_2, \ldots, \sigma_N\}$ where each σ_i takes the values ± 1 for $i = 1, \ldots, N$. We denote Σ_N the set of all σ. Then $|\Sigma_N| = 2^N$. Clearly Σ_N coincides with the space of all the possible 2^N Ising configurations of length N.

Definition 1 *Denote* $\{E_\sigma(N)\}_{\sigma \in \Sigma_N}$ *a family of* 2^N centered unit Gaussian *random variables:*

$$\mathrm{Av}\,(E_\sigma(N)) = 0\ , \tag{1}$$

and covariance matrix C_N *with elements defined by*

$$c_N(\sigma, \sigma) := \mathrm{Av}\left(E^2_\sigma(N)\right) = 1\ , \tag{2}$$

$$c_N(\sigma, \tau) := \mathrm{Av}\,(E_\sigma(N) E_\tau(N))\ . \tag{3}$$

Here $\mathrm{Av}\,(-)$ *denotes expectation with respect to the probability measure*

$$dP\,(E_1, \ldots, E_{2^N})\ =\ \frac{1}{\sqrt{(2\pi)^{2^N} det\,(C)}}\ e^{-\frac{1}{2}\langle E,\,C^{-1}E\rangle}\ dE_1 \cdots dE_{2^N}. \tag{4}$$

Definition 2

1. *For each N the Hamiltonian is given by*

$$H_N(\sigma) = -\sqrt{N} E_\sigma(N) \ . \tag{5}$$

2. *The partition function of the system is:*

$$Z_N(\beta, E) = \sum_\sigma e^{-\beta H_N(\sigma)} = \sum_\sigma e^{\beta\sqrt{N} E_\sigma(N)} \tag{6}$$

3. *The* quenched *free energy $f_N(\beta)$ of the system is defined as:*

$$-\beta f_N(\beta) := \alpha_N(\beta) := \frac{1}{N} \operatorname{Av}\left(\ln Z_N(\beta, E)\right) \ . \tag{7}$$

Remark 1 From now we write $E_\sigma(N) = E_\sigma$, dropping the N-dependence. Remark moreover that the above definition includes Gaussian families of the form

$$\begin{aligned} E_\sigma(N) \ = \ & J_0 + \sum_i J_i\sigma_i + \sum_{i,j} J_{i,j}\sigma_i\sigma_j + \sum_{i,j,k} J_{i,j,k}\sigma_i\sigma_j\sigma_k + \\ & + \ldots + \sum_{i_1,i_2,\ldots,i_N} J_{i_1,i_2,\ldots,i_N}\sigma_{i_1}\sigma_{i_2}\ldots\sigma_{i_N} \end{aligned} \tag{8}$$

in which every J is an indipendent Gaussian variable.

Examples.

1. The SK model. Consider first the model defined by

$$E_\sigma := \frac{1}{N} \sum_{i,j=1}^{N} J_{i,j}\sigma_i\sigma_j \tag{9}$$

where the $J_{i,j}$ are N^2 i.i.d. unit Gaussian random variables. A short computation yields

$$\mathrm{Av}(E_\sigma E_\tau) = [q_N(\sigma,\tau)]^2$$

where, as usual

$$q_N(\sigma,\tau) := \frac{1}{N}\sum_{k=1}^{N} \sigma_k \tau_k \tag{10}$$

is the overlap between the σ and τ spin configurations. The standard SK model is instead defined by

$$E^{SK}_\sigma := \frac{1}{N}\sum_{i<j=1}^{N} J_{i,j}\sigma_i\sigma_j \ . \tag{11}$$

However the quenched free energy densities (7) of the two models coincide up to a rescaling of the temperature, i.e.:

$$\alpha_N^{SK}(\sqrt{2}\beta) \ = \ \alpha_N(\beta)\,, \tag{12}$$

In fact, $J_{i,j}\sigma_i\sigma_j$ are centered, unit and i.i.d. Gaussian random variables $\forall\,(i,j)$, and $J_{i,j}\sigma_i\sigma_j = J_{j,i}\sigma_j\sigma_i$. Hence $J_{i,j}\sigma_i\sigma_j + J_{j,i}\sigma_j\sigma_i \stackrel{\mathcal{D}}{=} \sqrt{2}J_{i,j}\sigma_i\sigma_j$ (here $\stackrel{\mathcal{D}}{=}$ denotes equality in distribution of two random variables). Therefore, taking into account also the N diagonal terms:

$$\sqrt{N}E_\sigma \ \stackrel{\mathcal{D}}{=} \ \sqrt{N}\sqrt{2}E^{SK}_\sigma + J \ , \tag{13}$$

where J is a centered unit Gaussian variable. By (6,7) formula (13) immediately yields the relation (12).

2. The p-spin models. Here we consider the model:

$$E_\sigma := \sqrt{\frac{1}{N^p}}\sum_{i_1,\ldots,i_p=1}^{N} J_{i_1,\ldots,i_p}\sigma_{i_1}\cdots\sigma_{i_p} \tag{14}$$

where the $J_{i_1,\ldots,i_p}$ are once more i.i.d. unit Gaussian random variables. As before, a short computation yields

$$\mathrm{Av}(\mathrm{E}_\sigma \mathrm{E}_\tau) = [q_N(\sigma,\tau)]^p \tag{15}$$

3. The Derrida REM. Here the model is specified by Definition 1 with

$$\mathrm{Av}(E_\sigma E_\tau) = \delta(\sigma,\tau) \tag{16}$$

4. The Derrida-Gardner GREM. Its inclusion into the above framework is described in detail in Section 3.3.

Definition 3 *For each $\sigma \in \Sigma_N$ let π_1 and π_2 be the two canonical projections over the two subsets Σ_{N_1} and Σ_{N_2}, generated by a partition $\mathcal{P}$ of the coordinates $(\sigma_1,\ldots,\sigma_N)$ into a subset of N_1 coordinates and into a complementary set of N_2 coordinates: $N_1+N_2=N$, $\Sigma_N=\Sigma_{N_1}\times\Sigma_{N_2}$, $\pi_1\otimes\pi_2=1_{\Sigma_N}$.*

(Example: $N=4$; $\sigma\in\Sigma_4$ with coordinates denoted $\{\sigma_1,\sigma_2,\sigma_3,\sigma_4\}$. Consider for $N_1=N_2=2$ the partition $\mathcal{P}\sigma=(\sigma_1,\sigma_2)\cup(\sigma_3,\sigma_4)$. Then $\Sigma_N=\Sigma_{N_1}\times\Sigma_{N_2}$ and the two projections $\pi_k:\Sigma_N\to\Sigma_{N_k}$, $k=1,2$ act in the following way: $\pi_1(\sigma_1,\sigma_2,\sigma_3,\sigma_4)=(\sigma_1,\sigma_2)$ and $\pi_2(\sigma_1,\sigma_2,\sigma_3,\sigma_4)=(\sigma_3,\sigma_4)$).
Our main result is the following:

THEOREM 1 *Let the covariance matrices C_N fulfill the condition:*

$$c_N(\sigma,\tau) - \frac{N_1}{N}c_{N_1}(\pi_1(\sigma),\pi_1(\tau)) - \frac{N_2}{N}c_{N_2}(\pi_2(\sigma),\pi_2(\tau)) \le 0\ , \tag{17}$$

for every $N\ge\tilde N$, every $(\sigma,\tau)\in\Sigma_N\times\Sigma_N$ and every decomposition $N_1+N_2=N$. Then the thermodynamical limit exists, in the sense that

$$\lim_{N\to\infty}\frac{1}{N}\mathrm{Av}(\log Z_N(\beta)) = \sup_N \frac{1}{N}\mathrm{Av}(\log Z_N(\beta))\ . \tag{18}$$

Remark 2 The result (18) can be extended to the almost-everywhere convergence of free energy density, internal energy and ground state energy with elementary probability methods (see [GuTo])

Remark 3 The conditions (17) are not *necessary.* The proof itself will show that we only need the sign of the quantity in the left hand side of (17) in average, not pointwise. Moreover the condition (1) can be replaced by a more general small deviation vanishing for large N and (2) by a uniform (in N) bound over the diagonal terms. We plan to return over such a general case elsewhere.

Remark 4 It is still an open interesting question whether the class of models we control the thermodynamical limit of do have, in that limit, the properties axiomatically introduced by Ruelle in [Ru1] to define directly the infinite particle systems. To this purpose see [BS], [BoKu1], [BoKu2] and [BoKu3].

2 Proof

Within this section it is useful to consider 2 identical *copies* of the same system: the system 1 is assigned the Hamiltonian $H(\sigma)$ and the system 2 the Hamiltonian $H(\tau)$.

Definition 4 *The quenched measure over the two copies $\langle - \rangle$ is defined by*

$$< - >= \mathrm{Av}[Z(\beta,E)]^{-2} \sum_{(\sigma,\tau)\in\Sigma_N\times\Sigma_N} - \, e^{\beta(H(\sigma)+H(\tau))} \,. \qquad (19)$$

The definition may of course be generalized to r copies.

We want now to embed a Gaussian system $\{E_\sigma\}_{\Sigma_K}$ into a larger one $\{E_\tau\}_{\Sigma_L}$ for some $K < L$. In particular we want to embed two of them of size N_1 and N_2 into one of size $N = N_1 + N_2$. Our embedding procedure is defined in terms of the two canonical projections π_j, $j = 1, 2$ from Σ_N to Σ_{N_j} given in Definition (3).

Definition 5 *Given the family* $\{E_\mu\}_{\Sigma_{N_1}}$ *of size* N_1 *we* lift *it to one of size* N*:* $\{E^{(1)}_\sigma\}_{\Sigma_N}$ *defining*

$$E^{(1)}_\sigma \overset{\mathcal{D}}{=} E_{\pi_1(\sigma)} \ . \tag{20}$$

Moreover starting from $\{E_\mu\}_{\Sigma_{N_2}}$ *we define in the same way* $\{E^{(2)}_\sigma\}_{\Sigma_N}$ *by*

$$E^{(2)}_\sigma \overset{\mathcal{D}}{=} E_{\pi_2(\sigma)} \ . \tag{21}$$

Having defined each family $\{E_\sigma\}_{\Sigma_N}$, $\{E^{(1)}_\sigma\}_{\Sigma_{N_1}}$ *and* $\{E^{(2)}_\sigma\}_{\Sigma_{N_2}}$ *we specify their joint distribution requiring mutual independence.*

Remark 5 *:* The embedded Gaussian systems $\{E^{(1)}_\sigma\}_{\Sigma_{N_1}}$ and $\{E^{(2)}_\sigma\}_{\Sigma_{N_2}}$ are degenerate: In fact for all σ and τ such that $\pi_1(\sigma) = \pi_1(\tau)$

$$E^{(1)}_\sigma = E^{(1)}_\tau \ . \tag{22}$$

Summarizing we define the joint measure of $\{E_\sigma\}_{\Sigma_N}$, $\{E^{(1)}_\sigma\}_{\Sigma_{N_1}}$ and $\{E^{(2)}_\sigma\}_{\Sigma_{N_2}}$ $d\hat{P} = dPdP_1dP_2$ defined by the three covariances C_N, C_{N_1} and C_{N_2}.

Proof of THEOREM 1.

We proceed in three steps.

0) **Interpolation**

Given a pair (π_1, π_2) as before, following [GuTo], we pick three *independent* Gaussian systems $E^{(j)}_{\pi_j(\sigma)}$, $j = 0, 1, 2$ and introduce the quantity $(\pi_0(\sigma) = \sigma)$

$$H_{(N,N_1,N_2)}(\sigma,t) \;:=\; -\sum_{j=0}^{2} \sqrt{t_j N_j} E^{(j)}_{\pi_j(\sigma)} \tag{23}$$

where $t_0 = t$ and $t_1 = t_2 = (1-t)$, and the correspondent partition sum

$$Z_N(t,\beta) \;:=\; \sum_{\sigma\in\Sigma_N} e^{-\beta H_{(N,N_1,N_2)}(\sigma,t)}. \tag{24}$$

It is now easy to see that:

$$Z_N(1,\beta) = Z_N(\beta)\ , \tag{25}$$

and

$$\begin{aligned} Z_N(0,\beta) &= \sum_{\sigma\in\Sigma_N} e^{\beta(\sqrt{N_1}E^{(1)}_{\pi_1(\sigma)} + \sqrt{N_2}E^{(2)}_{\pi_2(\sigma)})} \\ &= \sum_{\tau\in\Sigma_{N_2}} \ \sum_{\sigma\in\Sigma_N;\, \pi_2(\sigma)=\tau} e^{\beta(\sqrt{N_1}E^{(1)}_{\pi_1(\sigma)} + \sqrt{N_2}E^{(2)}_{\tau})} \\ &= \sum_{\tau\in\Sigma_{N_2}} e^{\beta\sqrt{N_2}E^{(2)}_{\tau}} \sum_{\gamma\in\Sigma_{N_1}} e^{\beta\sqrt{N_1}E^{(1)}_{\gamma}} \\ &= Z_{N_1}(\beta)\cdot Z_{N_2}(\beta) \end{aligned} \tag{26}$$

1) **Boundedness**

The Jensen inequality

$$\mathrm{Av}\,(\log Z) \;\le\; \log(\mathrm{Av}\,(Z)) \tag{27}$$

implies

$$\frac{1}{N}\mathrm{Av}\,(\log Z_N(\beta)) \;\le\; \log(2) + \frac{\beta^2}{2} \tag{28}$$

because by (6) $\mathrm{Av}\,(Z) = 2e^{\beta^2/2}$ after performing the Gaussian integration.

2) **Monotonicity**

Taking the t derivative of the logarithm of (24) we get: (here we abbreviate $H_{N,N_1,N_2} = H$)

$$\frac{d}{dt}\log Z_N(t) \;=\; \frac{\beta}{Z_N(t)} \sum_{\sigma\in\Sigma_N} \left(\sum_{k=0}^{2} \epsilon_k \sqrt{\frac{N_k}{t_k}} E^{(k)}_{\pi_k(\sigma)}\, e^{-\beta H(\sigma,t)} \right) , \qquad (29)$$

where $\epsilon_0 = 1$ and $\epsilon_1 = \epsilon_2 = -1$.

We now use the integration by parts formula for correlated Gaussian variables $\{\xi_i\}$ with covariance $c_{i,j}$, which states

$$\mathrm{Av}\,(\xi_j \cdot f) \;=\; \mathrm{Av}\left(\sum_{k=1}^{n} c_{j,k} \cdot \frac{\partial f}{\partial \xi_k} \right) . \qquad (30)$$

This yields

$$\mathrm{Av}\left(\frac{1}{\beta}\frac{d}{dt}\log Z_N(t) \right) = \sum_{\sigma\in\Sigma_N} \sum_{k=0}^{2} \epsilon_k \sqrt{\frac{N_k}{t_k}}\, \mathrm{Av}\left(\frac{E^{(k)}_{\pi_k(\sigma)}\, e^{-\beta H}}{Z_N(t)} \right) \qquad (31)$$

$$= \sum_{\sigma\in\Sigma_N} \sum_{k=0}^{2} \epsilon_k \sqrt{\frac{N_k}{t_k}}\, \mathrm{Av}\left(\sum_{\tau_k\in\Sigma_{N_k}} c_{N_k}(\pi_k(\sigma),\tau_k) \cdot \frac{\partial}{\partial E^{(k)}_{\tau_k}} \frac{e^{-\beta H}}{Z_N(t)} \right)$$

Given now $\tau_k \in \Sigma_{N_k}$ fixed, we calculate

$$\frac{\partial}{\partial E^{(k)}_{\tau_k}} \frac{e^{-\beta H(\sigma,t)}}{Z_N(t)} = \beta \frac{\sqrt{N_k t_k}\; \delta^{\pi_k(\sigma)}_{\tau_k}\; e^{-\beta H(\sigma,t)} \cdot Z_N(t) - e^{-\beta H(\sigma,t)} \cdot \frac{\partial Z_N}{\partial E^{(k)}_{\tau_k}}}{Z_N^2(t)}$$

$$= \;\; \beta \frac{\sqrt{N_k t_k}\; \delta^{\pi_k(\sigma)}_{\tau_k} e^{-\beta H(\sigma,t)} \cdot Z_N(t) - \sqrt{N_k t_k}\; e^{-\beta H(\sigma,t)} \cdot \sum_{\xi\in\Sigma_N,\, \pi_k(\xi)=\tau_k} e^{-\beta H(\xi,t)}}{Z_N^2(t)}$$

The term with $k = 0$ in formula (31) is easy to calculate and we get:

$$N\beta \mathrm{Av}\left(\sum_{\sigma\in\Sigma_N} \sum_{\tau\in\Sigma_N} c_N(\sigma,\tau) \left[\delta^{\sigma}_{\tau} \frac{e^{-\beta H(\sigma,t)}}{Z_N} - \sum_{\xi\in\Sigma_N} \delta^{\tau}_{\xi} e^{-\beta(H(\xi,t)+H(\sigma,t))} \right] \right) \;=$$

$$= \quad N\beta \mathrm{Av}\left(\sum_{\sigma\in\Sigma_N} c_N(\sigma,\sigma)\cdot \frac{e^{-\beta H(\sigma,t)}}{Z_N} \;-\; \sum_{(\sigma,\tau)\in\Sigma_N\times\Sigma_N} c_N(\sigma,\tau)\, e^{-\beta(H(\tau,t)+H(\sigma,t))}\right)$$

$$= \quad N\beta\langle 1 \,-\, c_N(\sigma,\tau)\rangle_t \ , \tag{32}$$

where $< - >_t$ is the quenched measure with respect to the Hamiltonian (23).
In the same way for the term $k = 1$ (and similarly for $k = 2$) we obtain:

$$N_1\beta\mathrm{Av}\left(\sum_{\sigma\in\Sigma_N}\sum_{\tau\in\Sigma_{N_1}} c_{N_1}(\pi_1(\sigma),\tau)\left[\delta^{\tau}_{\pi_1(\sigma)}\frac{e^{-\beta H(\sigma,t)}}{Z_N} - \sum_{\xi\in\Sigma_N}\delta^{\tau}_{\pi_1(\xi)}e^{-\beta(H(\xi,t)+H(\sigma,t))}\right]\right) =$$

$$= N_1\langle 1 \,-\, c_{N_1}(\pi_1(\sigma),\pi_1(\tau))\rangle_t \ . \tag{33}$$

Summing up the three contributions we obtain:

$$\frac{1}{N}\frac{d}{dt}\mathrm{Av}\left(\log Z_N(t)\right) =$$

$$= -\beta^2 < c_N(\sigma,\tau) - \frac{N_1}{N}c_{N_1}(\pi_1(\sigma),\pi_1(\tau)) - \frac{N_2}{N}c_{N_2}(\pi_2(\sigma),\pi_2(\tau)) >_t \ , \tag{34}$$

and, by the hypothesis (17):

$$\frac{d}{dt}\mathrm{Av}\left(\log Z_N(t)\right) \geq 0 \ . \tag{35}$$

Formula (35) together with the boundary conditions (25) and (26) gives for every $N_1 + N_2 = N$

$$\alpha_N \geq \frac{N_1}{N}\alpha_{N_1} + \frac{N_2}{N}\alpha_{N_2} \ . \tag{36}$$

This entails Theorem 1 as explained for instance in [Ru2].

3 Examples

3.1 The SK and even p-spin models

For the sake of completeness we recover here the Guerra-Toninelli result [GuTo]. First note that by the definition (10) we have

$$q_N(\sigma,\tau) \,-\, \frac{N_1}{N} q_{N_1}(\pi_1(\sigma),\pi_1(\tau)) \,-\, \frac{N_2}{N} q_{N_2}(\pi_2(\sigma),\pi_2(\tau)) \;=\; 0\;. \tag{37}$$

so that (17) holds as an equality for $p = 1$ (the random field model). By (36) this means that the random field model free energy density doesn't depend on the size: $\alpha_N = \alpha_1$. For $p = 2u$ (SK corresponds to $u = 1$) formula (37) together with the convexity of the function $x \to x^{2u}$ implies (17):

$$q_N^{2u}(\sigma,\tau) \,-\, \frac{N_1}{N} q_{N_1}^{2u}(\pi_1(\sigma),\pi_1(\tau)) \,-\, \frac{N_2}{N} q_{N_2}^{2u}(\pi_2(\sigma),\pi_2(\tau)) \;\leq\; 0\;. \tag{38}$$

For the standard p-spin model defined as

$$E_\sigma = \sqrt{\frac{p!}{2N^p}} \sum_{i_1<\ldots<i_p} J_{i_1,\ldots,i_p}\sigma_{i_1}\cdots\sigma_{i_p} \tag{39}$$

we refer to [GuTo]

3.2 The REM

The model is defined by:

$$\mathrm{Av}\,(E_\sigma E_{\sigma'}) = \delta_{\sigma,\sigma'}\,. \tag{40}$$

Condition (17) is verified because it becomes

$$\delta_{\sigma,\sigma'} \leq \frac{N_1}{N}\delta_{\pi_1(\sigma),\pi_1(\sigma')} + \frac{N_2}{N}\delta_{\pi_2(\sigma),\pi_2(s')}\;. \tag{41}$$

In fact if $\sigma = \sigma'$ the previous formula is an identity. If $\sigma \neq \sigma'$ the left hand side is 0 but the right hand side is not always zero. Let us take for instance $\sigma = (+,+)$ and $\sigma' = (+,-)$, $\pi_1(+,+) = +$, $\pi_1(+-) = +$, $\pi_2(+,+) = +$, $\pi_2(+,-) = -$. In that case the left hand side is zero and the right hand side is 1/2.

3.3 The GREM

In order to show that our scheme includes the Derrida-Gardner GREM [DeGa] let first shortly recall its construction and add few observations. The GREM considers 2^N Gaussian random energies $H(\mu) = \sqrt{N} E_\mu$. Their covariance is given after the assignment of a *rooted tree* with n layers and 2^N leaves, $n < N$. The root α_1^N-*furcates*, the vertices at the end of the first layer α_2^N-*furcate* etc. up to the vertices at the end of the $n-1$ layer which α_n^N-*furcates* on the 2^N leaves. The topological constraint implies $\prod_{i=1}^n \alpha_i^N = 2^N$ and the prime number decomposition theorem imposes that $\alpha_i^N = 2^{k_i}$ with $k_1 + k_2 + ...k_n = N$. It is interesting and useful to associate to each leave μ a configuration of signs $\{\sigma_1, \sigma_2, ..., \sigma_N\}$. This can be done observing that the $\alpha_1^N = 2^{k_1}$ branches emerging from the root identify canonically the configurations of k_1 spins, the successive branches the configuration of k_2 spins and so on. We have in this way associated to each leaf either a path (the only one joining the root to it) or a spin configuration. The model is finally specified by the formula $E(\mu) = \sum_{i=1}^n \epsilon_i^{(\mu)}$ where the ϵ_i are thrown according to n Gaussians with a $\mathrm{Av}(\epsilon_i) = 0$ and $\mathrm{Av}[(\epsilon_i)^2] = a_i$: to each branch of the tree we associate an independent ϵ whose distribution (through its variance) depends only at which layer is the branch. Defining $v^{(l)} = \sum_{i=1}^{l-1} a_i\,,\quad (v^{(0)} = 0$ and $v^{(1)} = 1)$ it is immediate to prove that if two paths μ and ν merge at the level l we have $\mathrm{Av}(E_\mu E_\mu) = v^{(l)}$. For the Derrida-Gardner process over a tree $\mathcal{T}_{n,N}$ we will use the symbol $\{\mathcal{E}, \mathcal{T}_{n,N}\}$.

Our theorem on the existence of the thermodynamical limit can be used to show that the thermodynamical limit exist for the GREM in the sense that if $\{\mathcal{E}, \mathcal{T}_{n,N}\}$ is assigned for a given n and all $N > n$ we may show that its free energy density is decreasing (and bounded) in N. In order to do so, starting from a process $\{\mathcal{E}, \mathcal{T}_{n,N_1}\}$ we build the process $\{\mathcal{E}^{(1)}_{\pi_1}, \mathcal{T}_{n,N}\}$ with $N = N_1 + N_2$ in the following way: at each vertex of the tree $\mathcal{T}_{n,N_1}$ sitting on the layer i we increase the multiplicity of the *furcation* by a factor $(\alpha_i)^{N_2}$ carrying the same value $\epsilon^{(1)}_i$ to the newly introduced branches. The new process will enjoy the property

$$\mathrm{Av}(E^{(1)}_{\pi_1(\sigma)} E^{(1)}_{\pi_1(\tau)}) \geq v^{(l)}\,. \tag{42}$$

We apply the same construction to build $\{\mathcal{E}^{(2)}_{\pi_2}, \mathcal{T}_{n,N}\}$ and we have

$$\mathrm{Av}(E^{(2)}_{\pi_2(\sigma)} E^{(2)}_{\pi_2(\tau)}) \geq v^{(l)}\,. \tag{43}$$

It is now straightforward to verify that conditions (42) and (43) imply (17).

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